STUDY OF LASER PLASMA INTERACTIONS USING AN EULERIAN VLASOV CODE

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Abstract

We use a one-dimensional Eulerian Vlasov Code to study laser-plasma interactions. The excitation of parametric instabilities, such as stimulated Raman scattering (SRS) and stimulated Brillouin scattering (SBS) are an important issue for the efficiency of inertial confinement fusion. Our aim is to understand the growth, saturation, and coupling of these instabilities, in particular the kinetic aspects of these phenomena.

The code evolves both species with the fully relativistic Vlasov equation in the direction of laser propagation. It also incorporates a cold fluid transverse velocity for each species, which is needed to couple the transverse and longitudinal dynamics. The code uses the method of fractional steps and cubic spline interpolation to advance the distribution function in time.

The plasma is finite and not periodic. On the boundaries we place "absorbing plates" that collect the particles flowing onto them. This charge enters Poisson's equation as a boundary condition on the electric field.

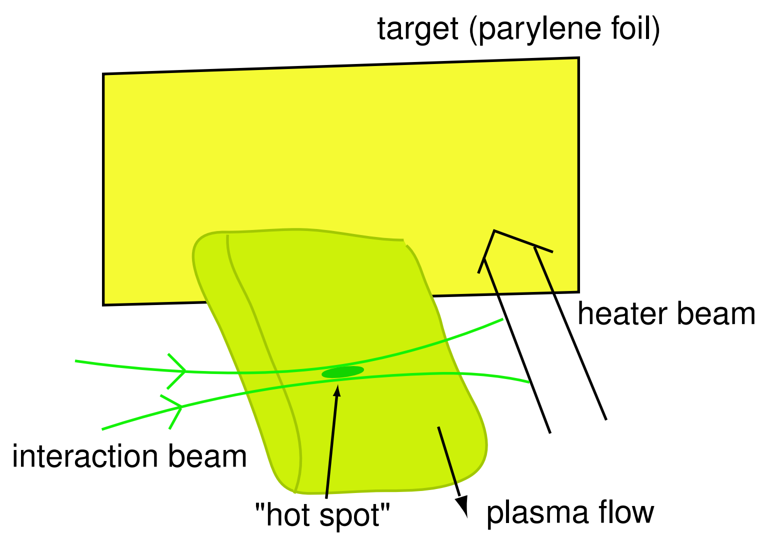
We are able to run the code for many times longer than it takes the laser to cross the computational domain. Once the laser passes a region of plasma, it modifies the equilibrium. This modification is sufficient for SRS to grow from; we do not "seed" SRS by imposing any perturbations or a secondary laser.

For parameters similar to the recent Trident single hot spot experiments, SRS is not limited by the Langmuir Decay Instability (LDI) or other ion processes, but by kinetic effects. The code clearly resolves the formation of large vortices in phase-space due to electron trapping in the field of the SRS electron plasma wave. We also see SBS starting to develop, but we have not run for long enough times for the ions to move through several ion-acoustic wave periods.

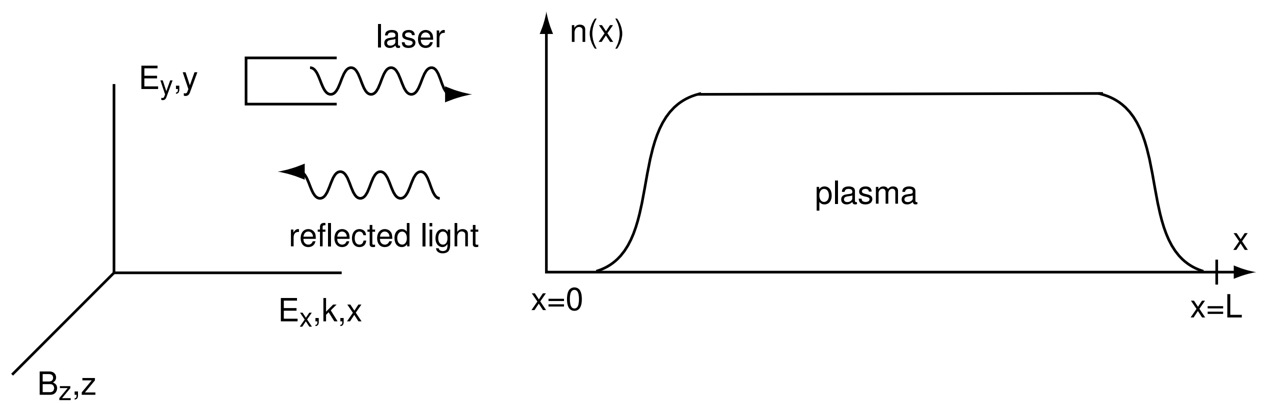
Work is underway to extend the code to $1\frac{1}{2}$ dimensions and to parallelize it.

Single Hot Spot Experiments

[D. S. Montgomery, J. A. Cobble, et al. *Phys. Plasmas*, 9(5):2311–2320, 2002.]



One-Dimensional Model Geometry



1-Dimensional Kinetic Model

- Spatial variation only in x: $\nabla \rightarrow \partial/\partial x$
- Fields: $\vec{E} = (E_x, E_y, 0)$ $\vec{B} = (0, 0, B_z)$
- Transverse cold beam: $F_s(x, p_x, P_y, t) = \delta(P_y) f_s(x, p_x, t)$, $P_y = m_s v_{ys} + q_s A_y = y$ canonical mom. y momentum equation becomes [Gabor, *Proc. IRE* 33:792, 1945]

$$\boxed{\frac{\partial v_{ys}}{\partial t} = \frac{q_s}{m_s} E_y}$$

- Longitudinal dynamics: Relativistic 1-D Vlasov Eqn. for $f_s(x, p_x, t)$

$$\boxed{\frac{\partial f_s}{\partial t} + v_x \frac{\partial f_s}{\partial x} + \frac{q_s}{m_s}(E_x + v_{ys} B_z) \frac{\partial f_s}{\partial p_x} = 0}$$

- Longitudinal field: E_x (electrostatic perturbations)

$$\text{Poisson} \longrightarrow \qquad \boxed{\frac{\partial E_x}{\partial x} = \frac{e}{\epsilon_0}(n_i - n_e)}$$

- Transverse fields: laser, electromagnetic perturbations

$$\text{Ampère} \longrightarrow \qquad \frac{\partial E_y}{\partial t} + c^2 \frac{\partial B_z}{\partial x} = -\frac{1}{\epsilon_0} J_y$$

$$\boxed{E^{\pm} \equiv E_y \pm c B_z \qquad \longrightarrow \qquad \left(\frac{\partial}{\partial t} \pm c \frac{\partial}{\partial x} \right) E^{\pm} = -\frac{1}{\epsilon_0} J_y}$$

The Time-Stepping Algorithm

Operator splitting with cubic spline for Vlasov; leapfrog for electromagnetic fields [Ghizzo, Bertrand, Shoucri *et al.* , *Journ. Comp. Phys.* 90(431), 1990]

t/dt	0	0^+	$\frac{1}{2}^-$	$\frac{1}{2}$	$\frac{1}{2}^+$	1^-	1	1^+
(leapfrog) f	$f_0 \xrightarrow{(\partial_t+v_x\partial_x)f=0}$		$f_{\frac{1}{2}^-}$	$\rightsquigarrow$ $(\partial_t+F_{1/2}\partial_{p_x})f=0, \quad F=q(E_x+v_yB_z)$	$f_{\frac{1}{2}^+} \xrightarrow{(\partial_t+v_x\partial_x)f=0}$		f_1	
(what we do) f	$f_0 \xrightarrow{(\partial_t+v_x\partial_x)f=0}$					$f_{1^-} \rightsquigarrow$ $(\partial_t+F_{1/2}\partial_{p_x})f=0$	f_1	
n	n_0			$n_{\frac{1}{2}} = \frac{1}{2}(n_0+n_1)$			$n_1[f_{1^-}]$	
N_p	N_{p0}			$N_{p\frac{1}{2}} = \frac{1}{2}(N_{p0}+N_{p1})$			$N_{p1}[n_1]$	
dN_w	dN_{w0}			$dN_{w\frac{1}{2}} = \frac{1}{2}(dN_{w0}+dN_{w1})$			$dN_{w1}[f_{1^-}]$	
N_w	N_{w0}			$N_{w\frac{1}{2}} = N_{w0} + \frac{1}{2}dN_{w\frac{1}{2}}$			$N_{w1} = N_{w0} + dN_{w1}$	
E_x				$E_{x\frac{1}{2}}[n_{\frac{1}{2}}, N_{p\frac{1}{2}}, N_{w\frac{1}{2}}]$				
$E^\pm$		$E^\pm_{0^+} \xrightarrow{(\partial_t\pm\partial_x)E^\pm=0}$		$E^\pm_{\frac{1}{2}} = \frac{1}{2}(E^\pm_0 + E^\pm_{1^-})$		$E^\pm_{1^-} \rightsquigarrow$ $(\partial_t\pm\partial_x)E^\pm=-\epsilon_0^{-1}J_{y1}$		$E^\pm_{1^+}$
v_y	$v_{y0} \xrightarrow{m\partial_t v_y=qE_{y\frac{1}{2}}}$			$v_{y\frac{1}{2}} = \frac{1}{2}(v_{y0}+v_{y1})$			v_{y1}	

Boundary Conditions

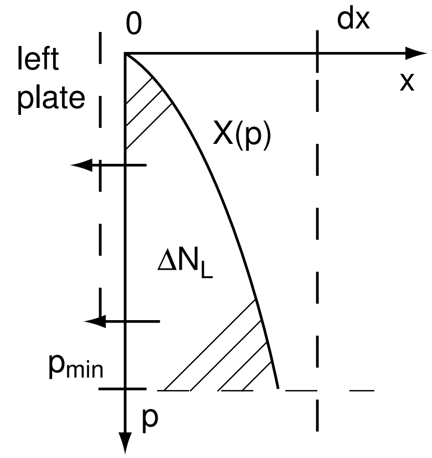
- Nonperiodic grids in x and p_x; plasma is finite.
- "Absorbing plates" (transparent to electromagnetic fields) at $x = 0$ and $x = L$. Particles that flow out to the plates stay there and no longer evolve.

$$N_{sL},\ N_{sP},\ N_{sR} = \#\text{ of particles of species } s \text{ on left plate, in plasma, on right plate}$$

$$X(p) = -dt\ v(p) \qquad v(p) = \frac{p}{\gamma m}$$

$$\Delta N_L = \int_{p_{min}}^{0} dp \int_0^{X(p)} dx f(x,p)$$

$$\approx -dt \int_{p_{min}}^{0} dp\ v\ f(0,p)$$



- Charges on plates enter E_x as a boundary condition.

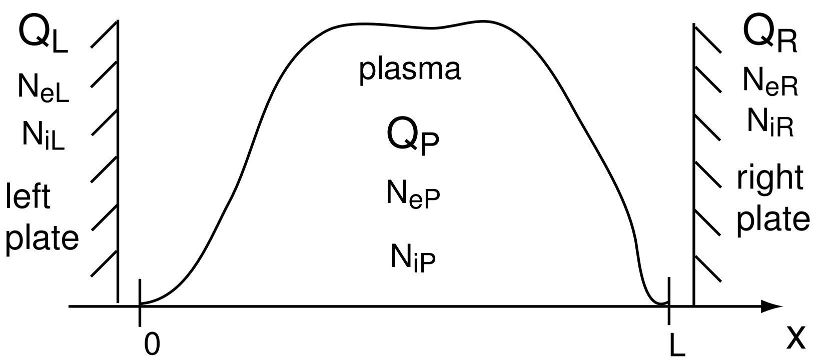
$$\frac{\partial E_x}{\partial x} = \epsilon_0^{-1}\rho$$

$$Q_L + Q_P + Q_R = 0$$

$$\therefore E_x(x < \text{left plate}) = 0$$

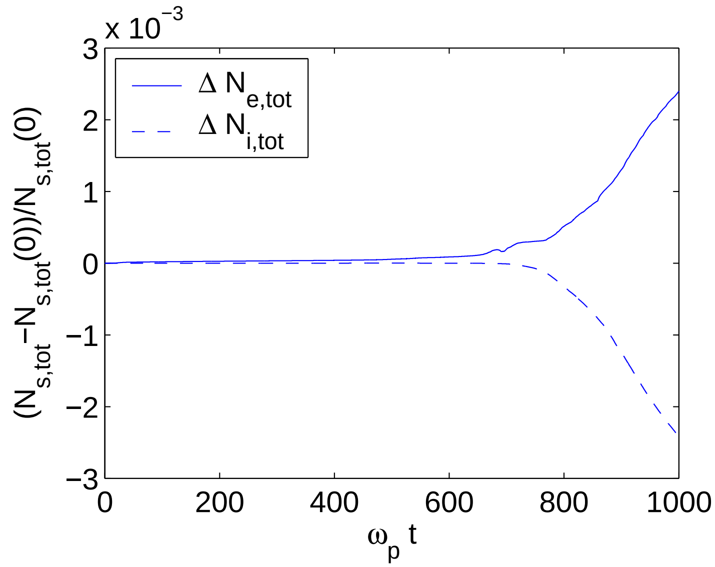
$$E(x=0) = \frac{1}{2\epsilon_0}(Q_L - Q_P - Q_R)$$

$$E(x) = E(x=0) + \epsilon_0^{-1}\int_0^x \rho(x')dx'$$

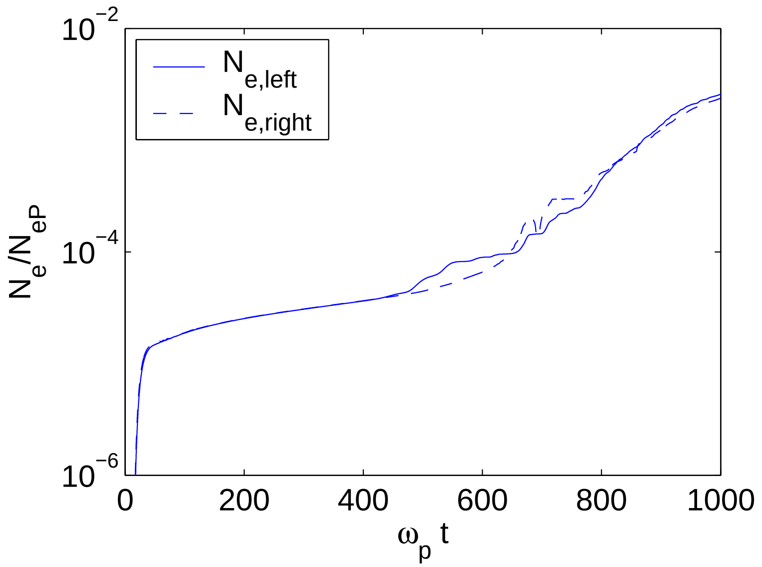


Outflow to plates

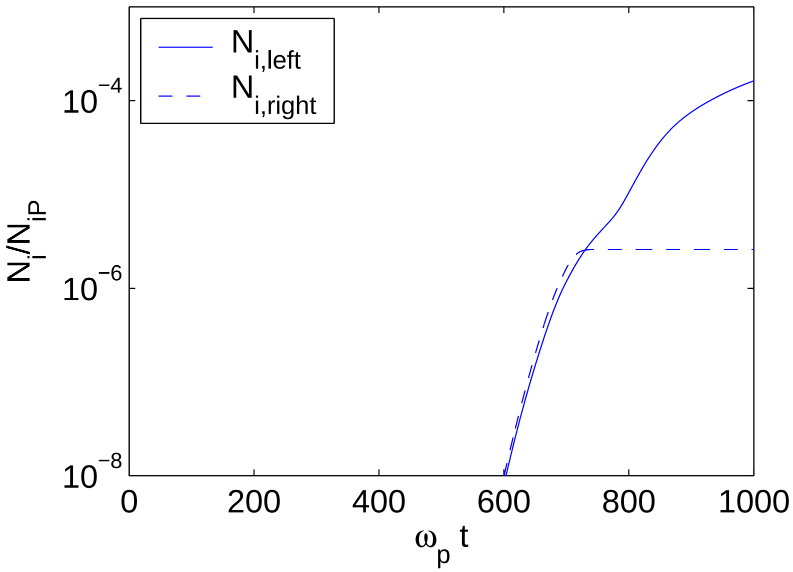
Change in total number of particles



Electrons on left, right walls



Ions on left, right walls



Run Parameters

- Scales used in code:

  time: $t_0 = 1/\omega_p$ length: $d_e = c/\omega_p$ speed: c momentum (all species): $p_0 = m_e c$

- Plasma Parameters:

  Plasma length: $L = 75\ \mu\text{m} = 160\ d_e$
  Plasma density: $n_e = 1.28 \cdot 10^{20}\ \text{cm}^{-3} = 0.032\ n_{crit}$
  initial electron Temperature: $T_e = 350$ eV initial ion Temperature: $T_i = 100$ eV
  electron thermal speed: $v_{Te} = 0.0262\ c$ ion thermal speed: $v_{Ti} = 3.27 \cdot 10^{-4}\ c$
  Plasma frequency: $\omega_p = 6.4 \cdot 10^{14}$ rad/s $\rightarrow$ $t_0 = 1.56 \cdot 10^{-15}$ s, $d_e = 470$ nm
  Debye length: $\lambda_{De} = 12.3\ \text{nm} = 0.0262\ d_e$

- Laser Parameters:

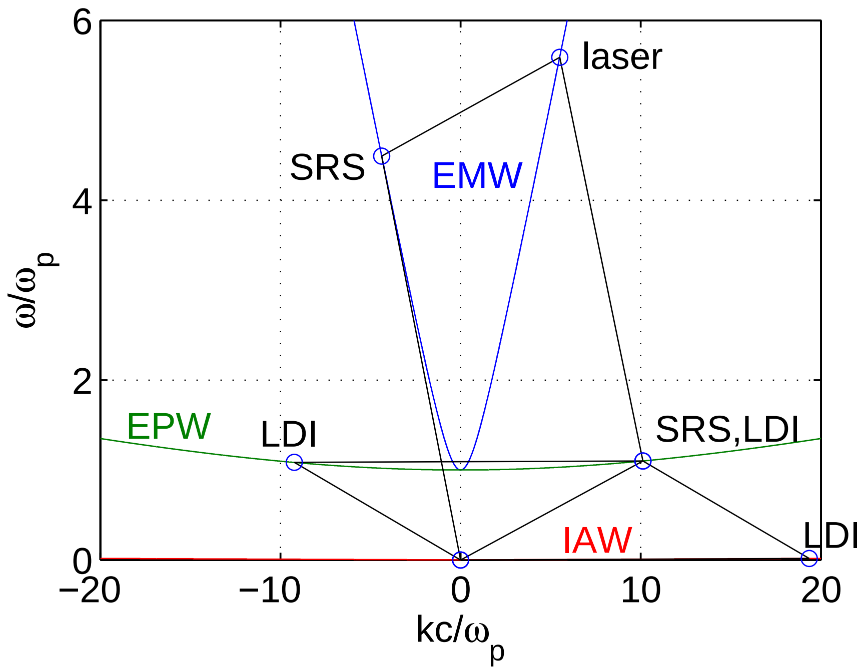
  Laser frequency: $\omega_0 = 3.57 \cdot 10^{15}\ \text{rad/s} = 5.59\ \omega_p$ free-space wavelength: $\lambda_0 = 527\ \text{nm} = 1.12\ d_e$
  Laser instensity: $I_0 = 10^{16}\ \text{W/cm}^2$
  electron quiver velocity: $v_{quiv} = .045\ c$

- Numerical parameters:

  Timestep: $dt = c\ dx = 0.01/\omega_p$
  Total runtime: $T = 1.56\ \text{ps} = 1000\ t_0$
  Spatial gridpoints: 16001 $\rightarrow dx = 0.01\ d_e$
  electron momentum gridpoints: 301 $\rightarrow dp_e = 0.0027\ m_e c = 0.102\ p_{Te}$
  ion momentum gridpoints: 151 $\rightarrow dp_i = 0.069\ m_e c = 0.187\ p_{Ti}$

Parametric Instabilities: Coupling of Modes

EPW (e^- plasma)	$\omega^2 = \omega_{pe}^2 + 3v_{Te}^2 k^2$	e/s $\quad \vec{E} \Vert \vec{k}$
IAW (ion acoustic)	$\omega^2 = c_a^2 k^2$	e/s $\quad \vec{E} \Vert \vec{k}$
EMW (electromagnetic)	$\omega^2 = \omega_{pe}^2 + c^2 k^2$	e/m $\quad \vec{E} \perp \vec{k}$



Matching: $\omega_1 = \omega_2 + \omega_3 \qquad k_1 = k_2 + k_3$

$$\vec{E}_j = a_j(x,t)\sqrt{\frac{2\omega_j}{\epsilon_0}}\hat{\epsilon}_j e^{(k_j x - \omega_j t)}$$

$$|a_j|^2 = \frac{\epsilon_0}{2\omega_j}|E_j|^2 = \text{action density}$$

$$\frac{Da_1}{D_1 t} = -K a_2 a_3$$

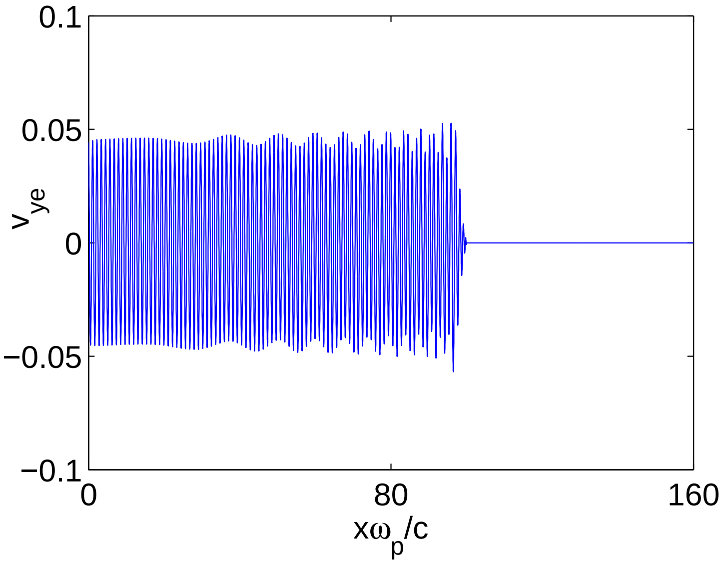
$$\frac{Da_2}{D_2 t} = K^* a_1 a_3^*$$

$$\frac{Da_3}{D_3 t} = K^* a_1 a_2^*$$

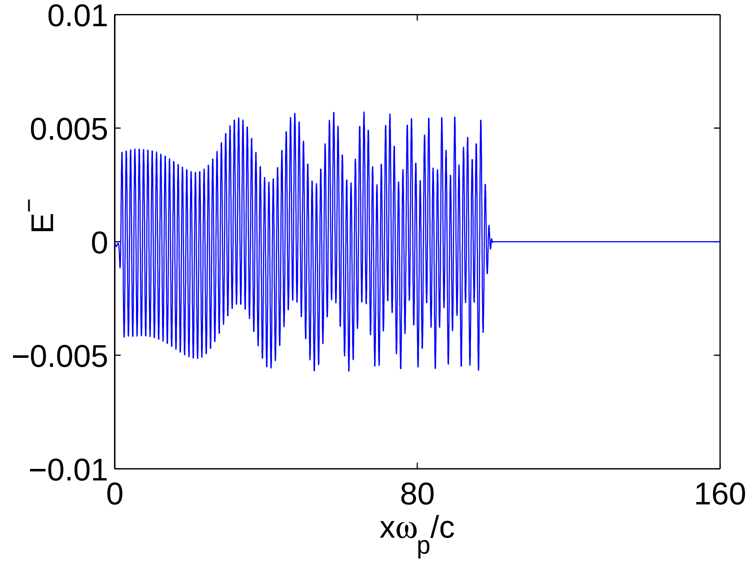
$$\frac{D}{D_j t} = \frac{\partial}{\partial t} + v_{gj}\frac{\partial}{\partial x} + \nu_j$$

Laser-Modified Equilibrium: Electromagnetic Fields

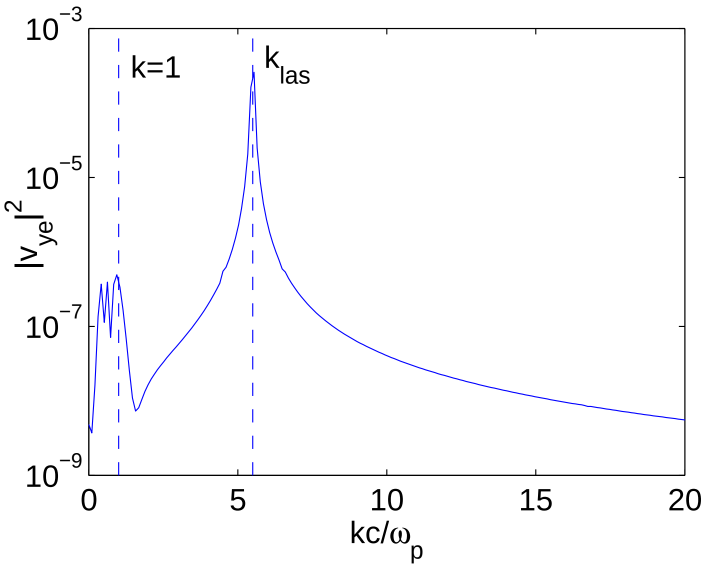
v_{ye} at $\omega_p t = 100$



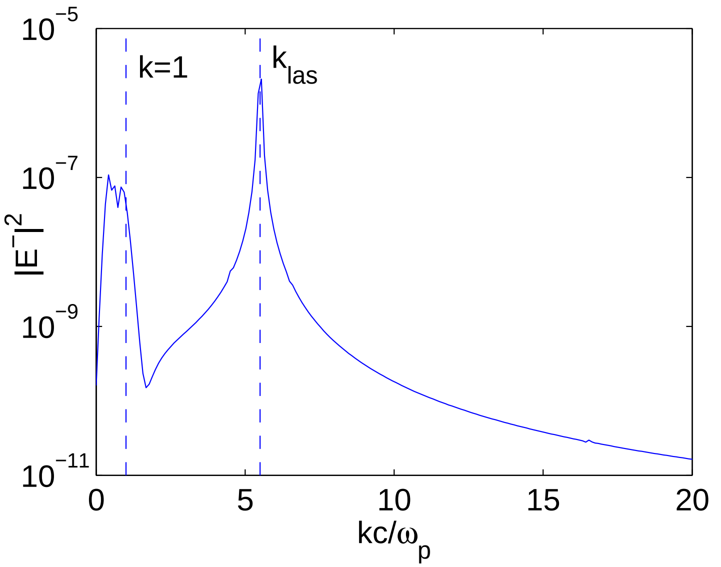
E^- at $\omega_p t = 100$



$|v_{ye}(\omega)|^2$ at $\omega_p t = 100, x\omega_p/c = 20 - 80$

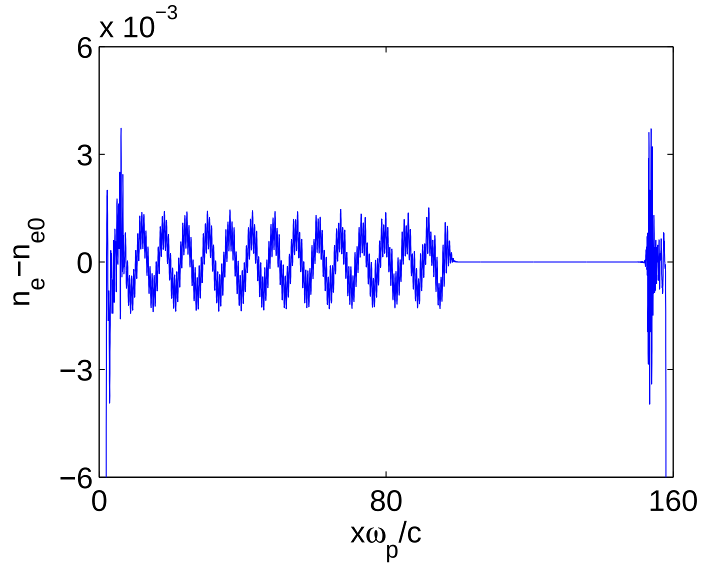


$|E^-(\omega)|^2$ at $\omega_p t = 100, x\omega_p/c = 20 - 80$

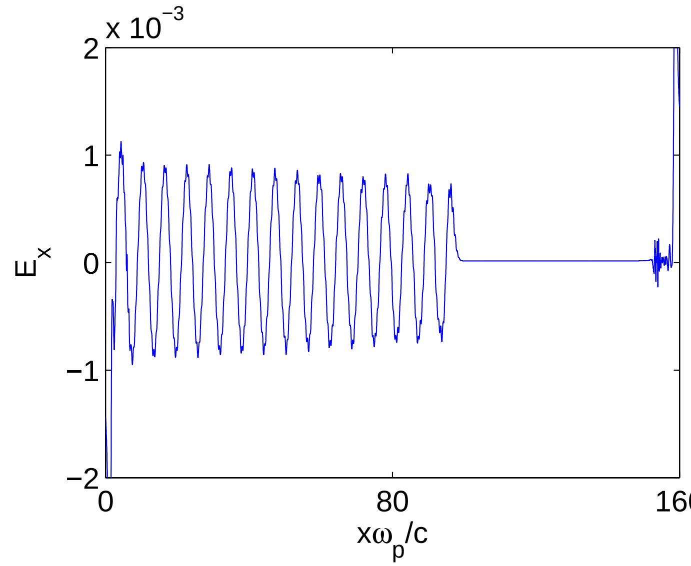


Laser-Modified Equilibrium: Electrostatic Fields

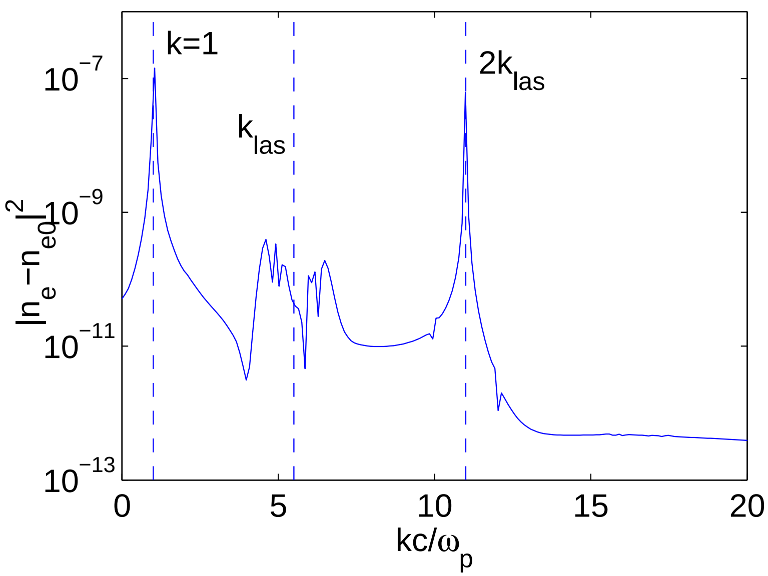
$n_e - n_{e0}$ at $\omega_p t = 100$



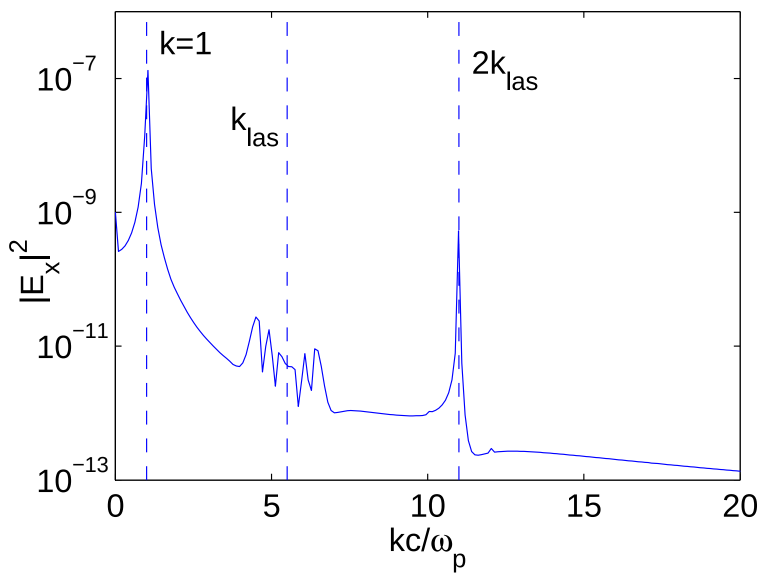
E_x at $\omega_p t = 100$



$|n_e(\omega)|^2$ at $\omega_p t = 100, x\omega_p/c = 20 - 80$

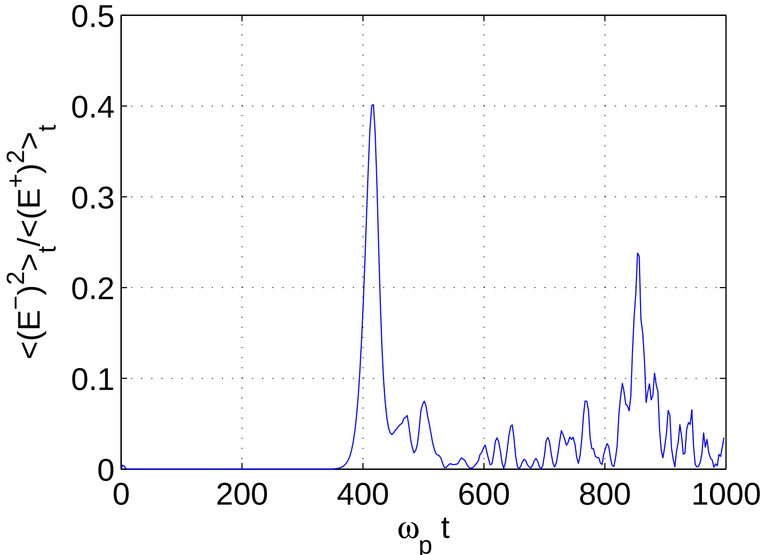


$|E_x(\omega)|^2$ at $\omega_p t = 100, x\omega_p/c = 20 - 80$

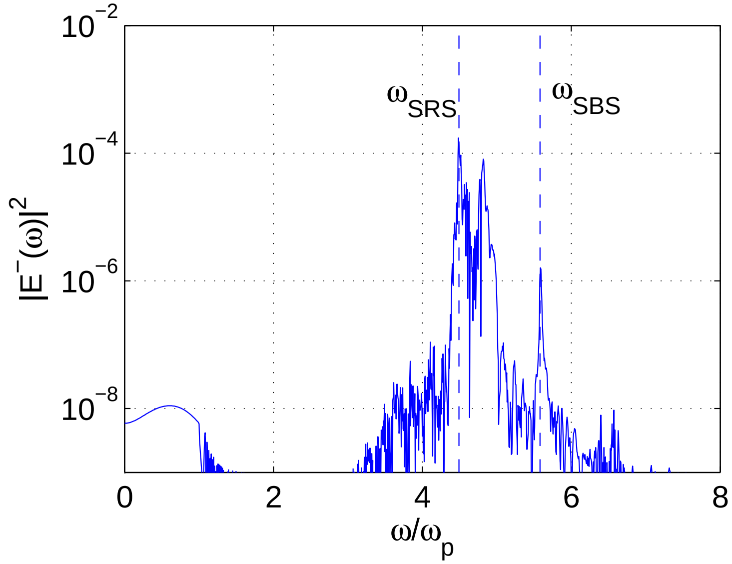


Reflected Light at $x = 0$

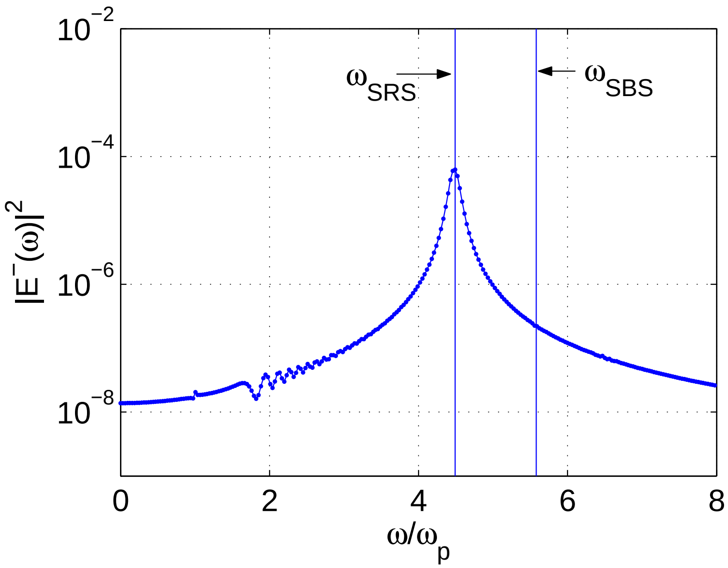
Instantaneous Reflectivity (Avg.=3.1%)



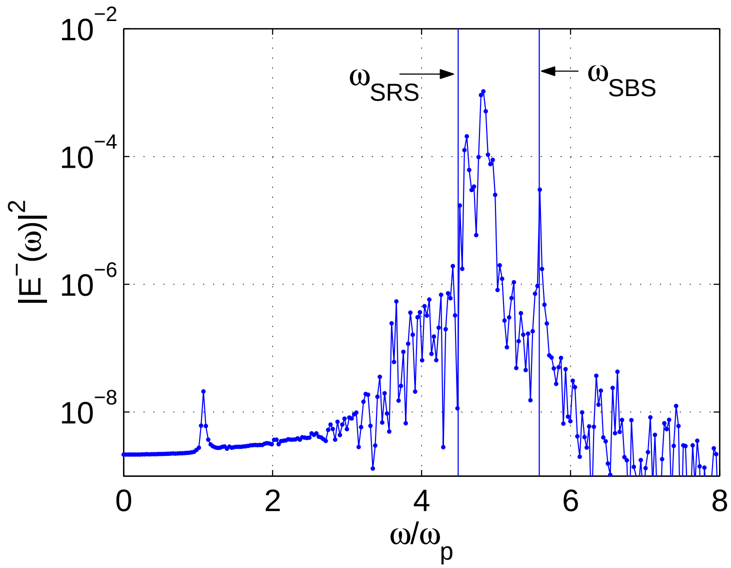
$|E^-(\omega)^2|$ at $x = 0$, $t\omega_p = 0 - 1000$



$|E^-(\omega)^2|$ at $x = 0$, $t\omega_p = 200 - 400$

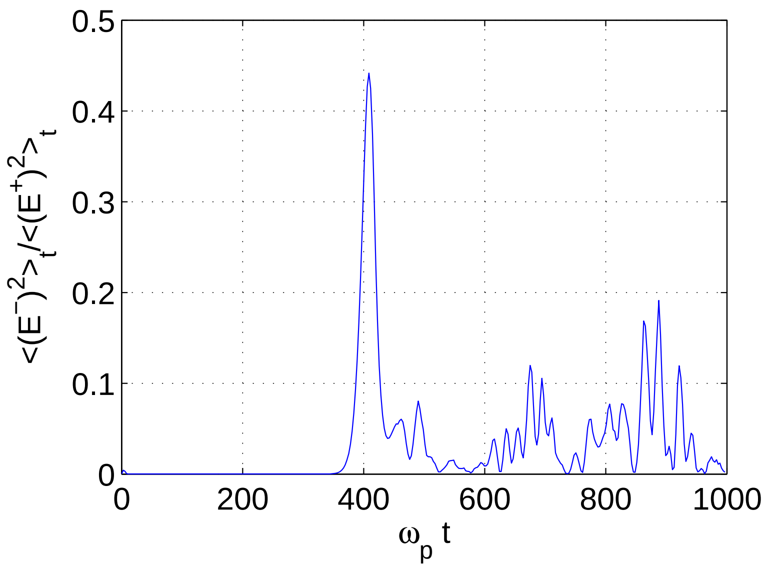


$|E^-(\omega)^2|$ at $x = 0$, $t\omega_p = 800 - 1000$

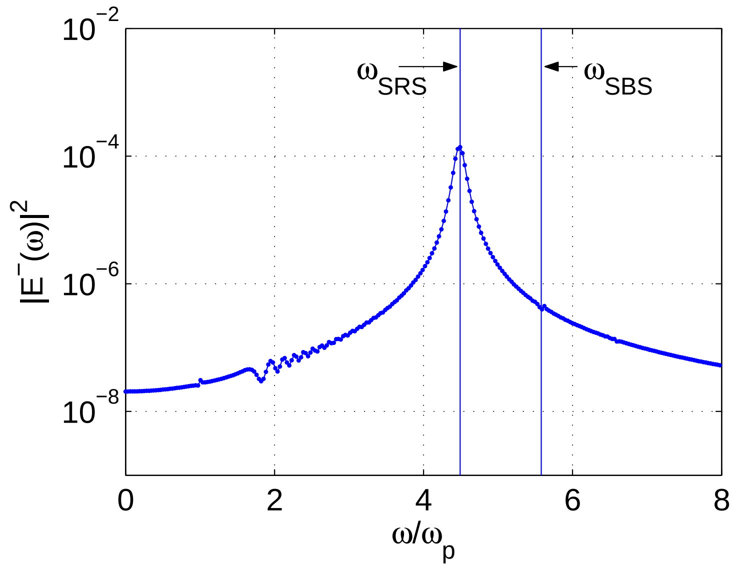


Reflected Light at $x = 0$, Stationary Ions

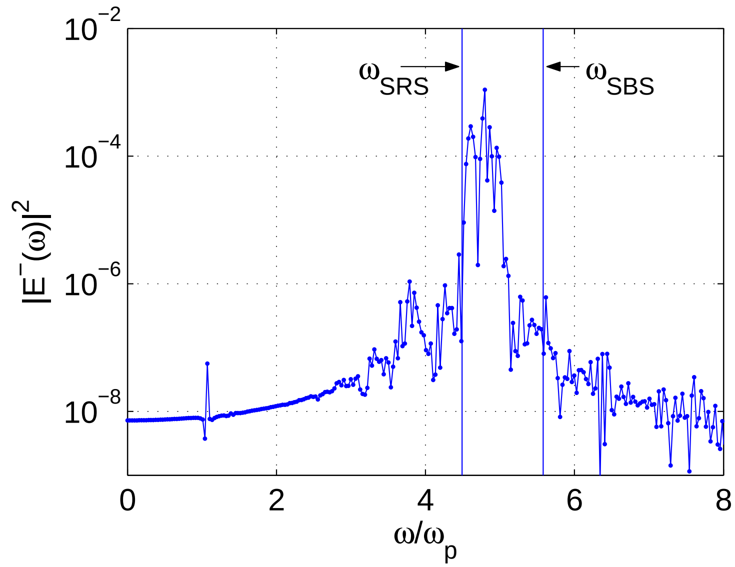
Instantaneous Reflectivity (Avg.: 3.4%)



$|E^-(\omega)^2|$ at $x = 0$, $t\omega_p = 200 - 400$

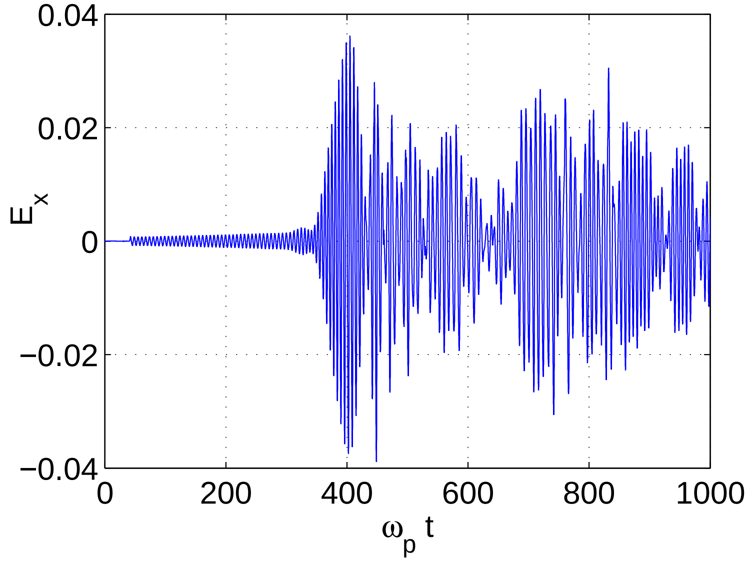


$|E^-(\omega)^2|$ at $x = 0$, $t\omega_p = 800 - 1000$

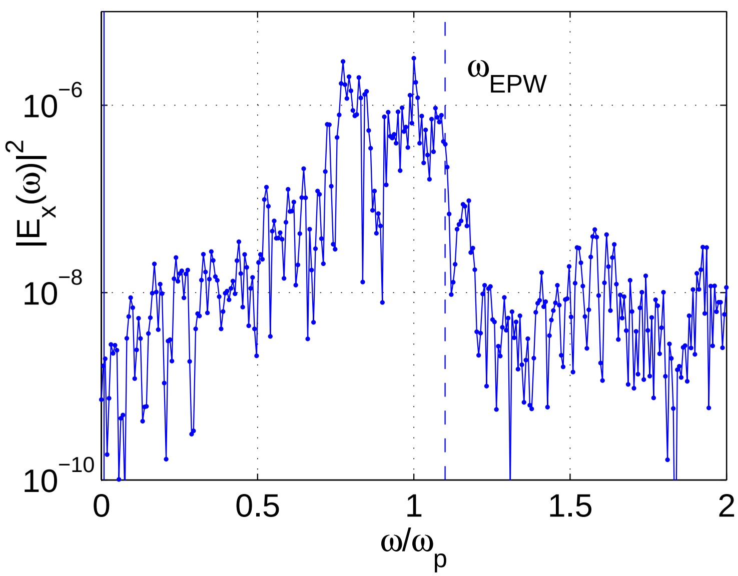


Longitudinal Electric Field E_x vs. t

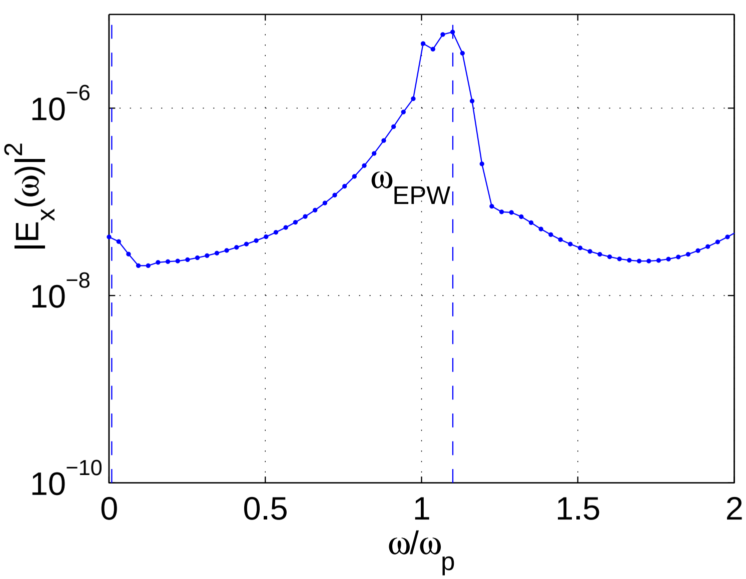
E_x at $x\omega_p/c = 40$



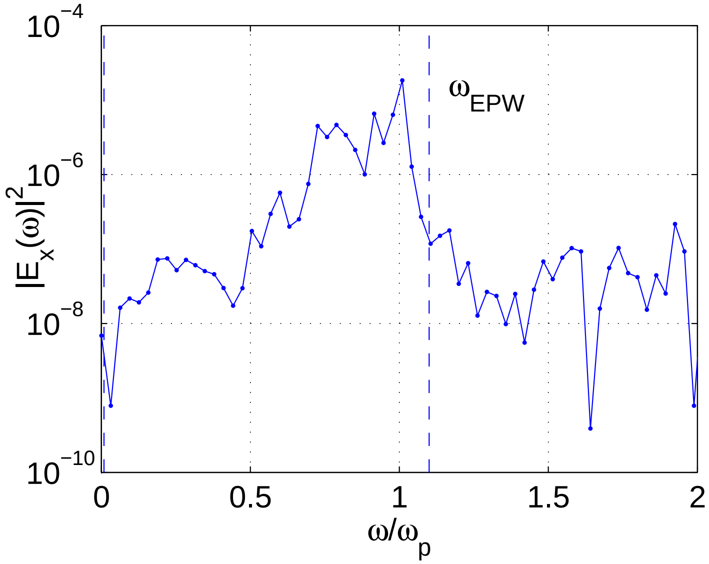
$|E_x(\omega)|^2$ at $x\omega_p/c = 40, \omega_p t = 0 - 1000$



$|E_x(\omega)|^2$ at $x\omega_p/c = 40, \omega_p t = 200 - 400$

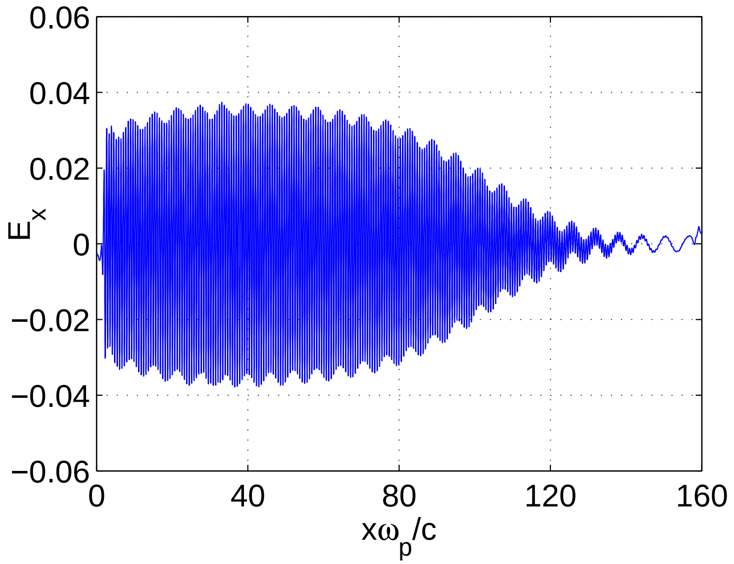


$|E_x(\omega)|^2$ at $x\omega_p/c = 40, \omega_p t = 800 - 1000$

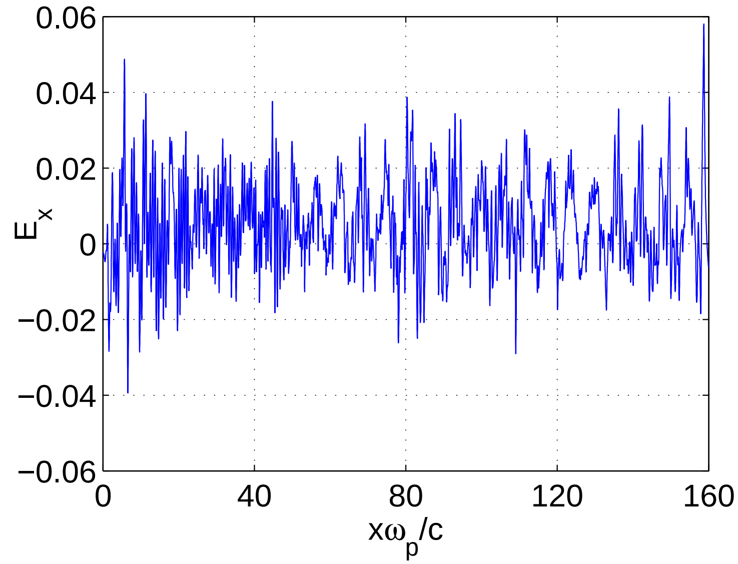


Longitudinal Electric Field E_x vs. x

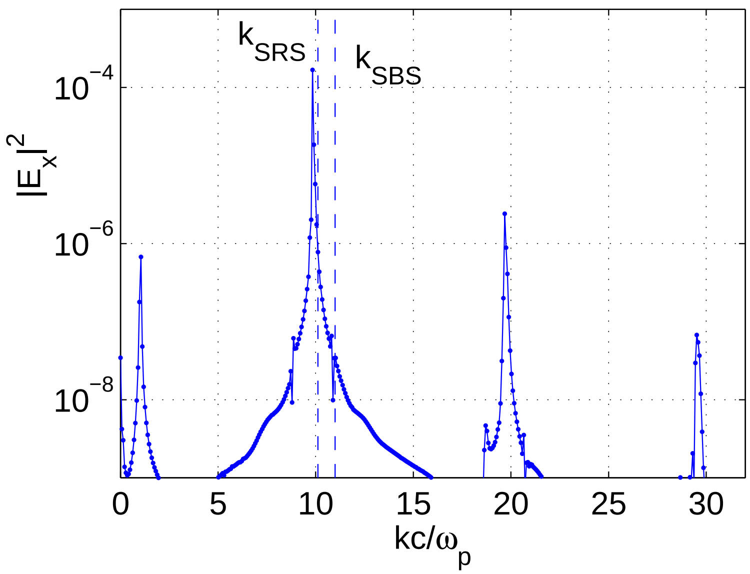
E_x at $\omega_p t = 400$



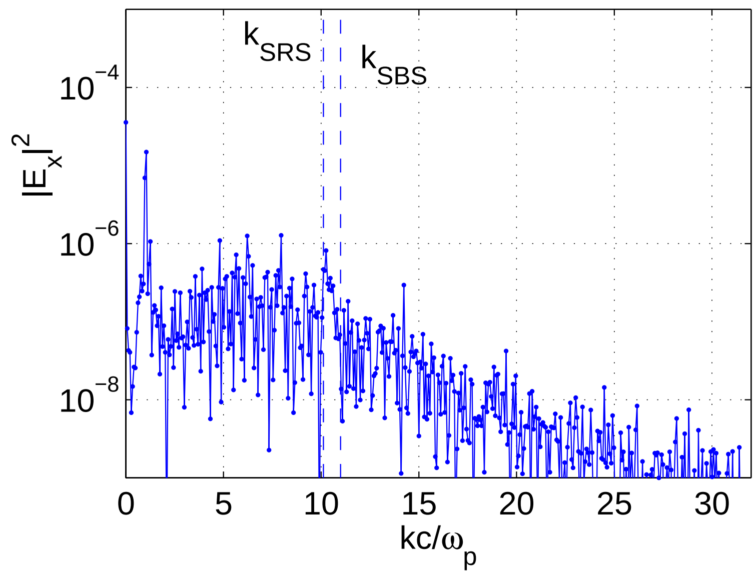
E_x at $\omega_p t = 900$



$|E_x(k)|^2$ at $\omega_p t = 400, x\omega_p/c = 30 - 120$

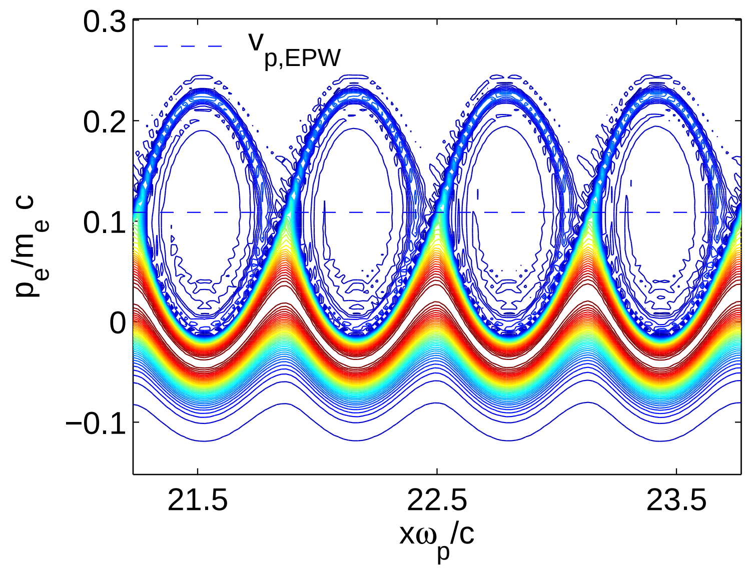


$|E_x(k)|^2$ at $\omega_p t = 900, x\omega_p/c = 30 - 120$

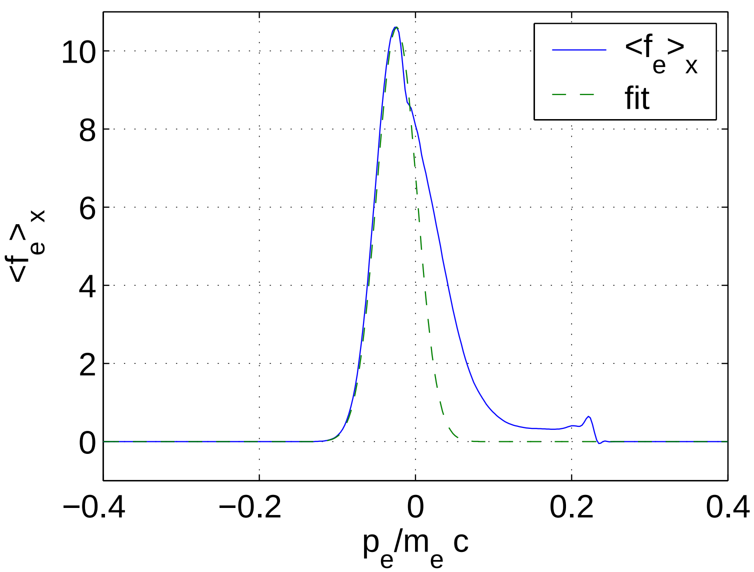


Trapped f_e during first SRS burst

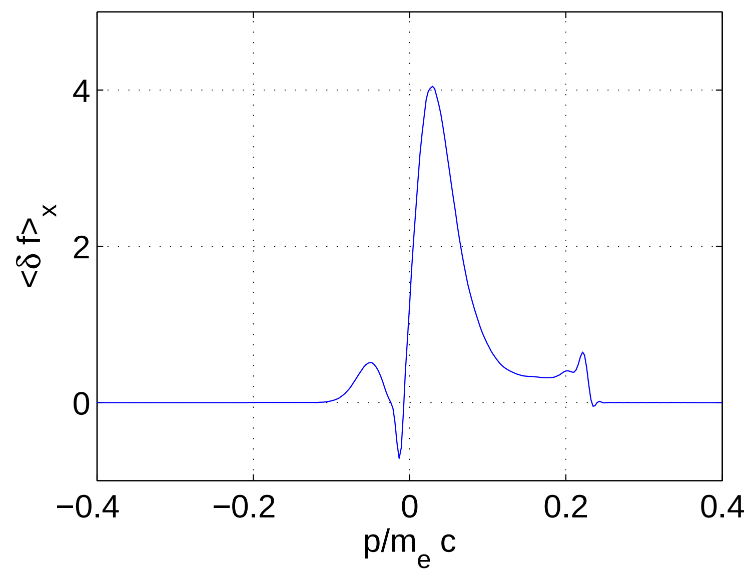
Contour Plot of f_e at $\omega_p t = 400$



Spatially-averaged $< f_e >_x$

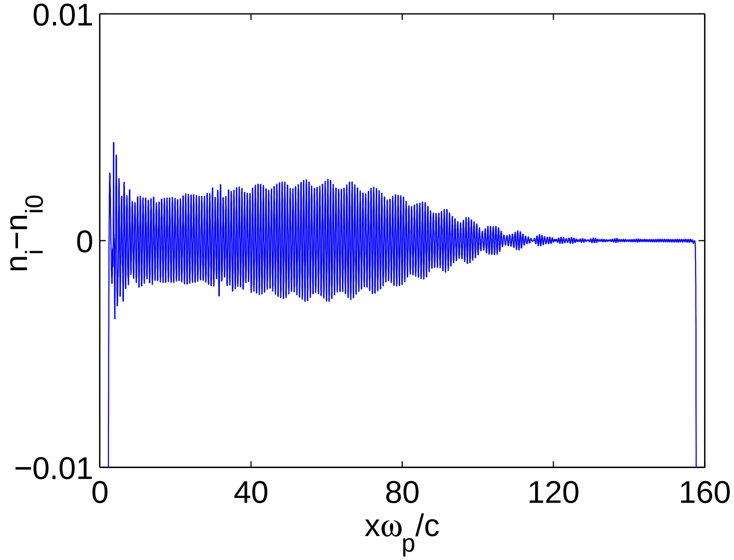


$< f_e >_x$ − Maxwellian fit

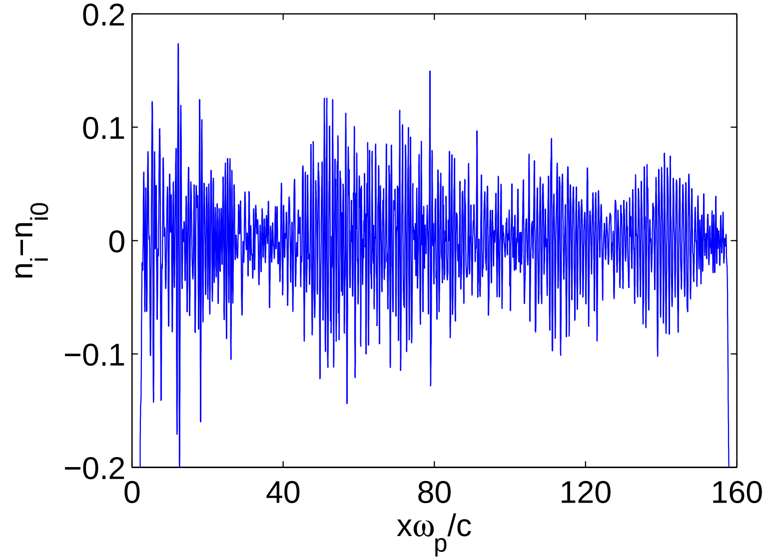


Ion Density Fluctuations vs. x

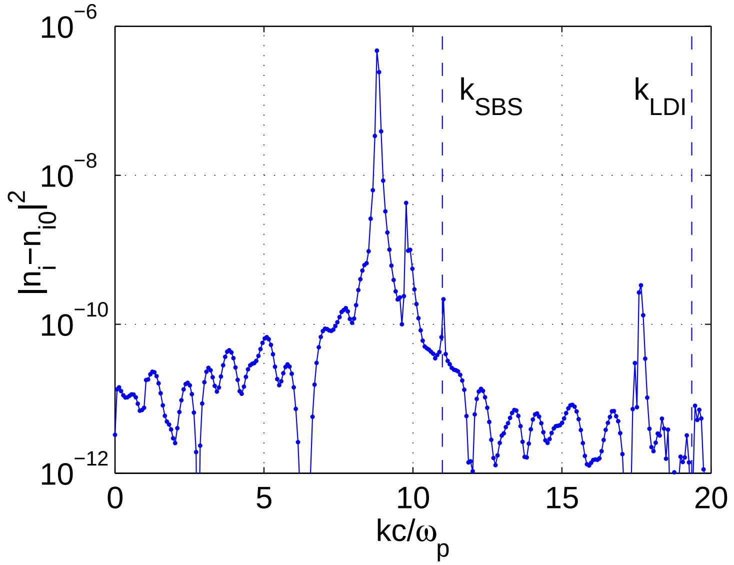
$n_i - n_{i0}$ at $\omega_p t = 450$



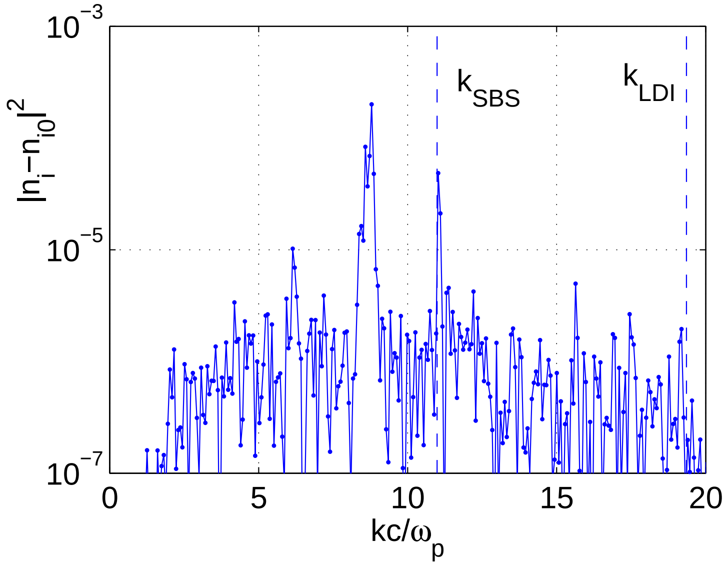
$n_i - n_{i0}$ at $\omega_p t = 1000$



$|n_i(\omega)|^2$ at $\omega_p t = 450, x\omega_p/c = 30 - 120$

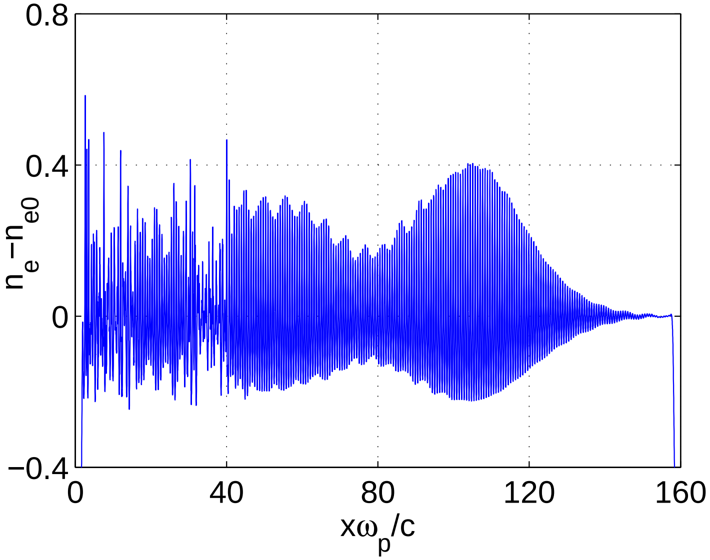


$|n_i(\omega)|^2$ at $\omega_p t = 1000, x\omega_p/c = 30 - 120$

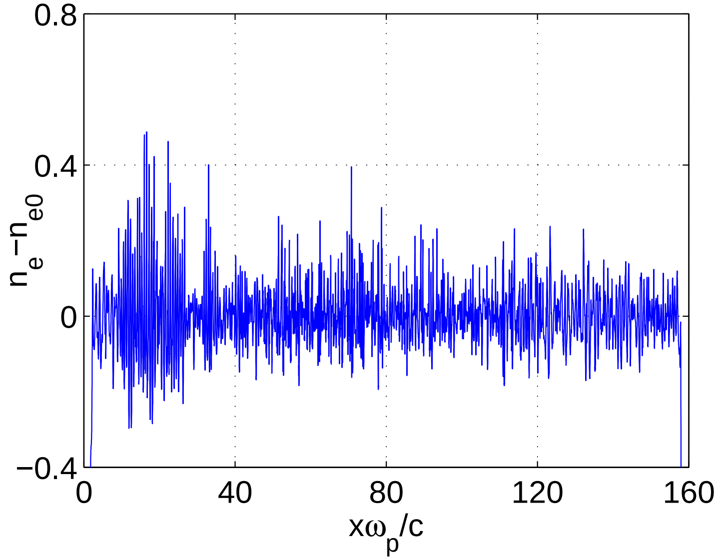


Electron Density Fluctuations vs. x

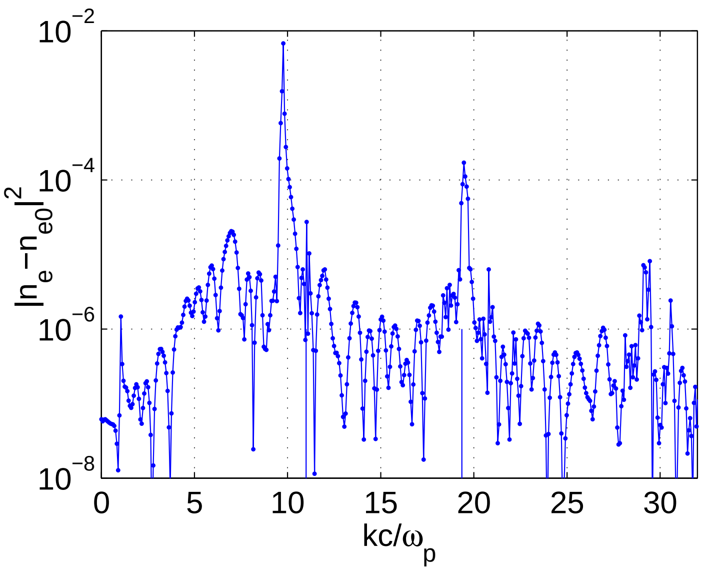
$n_e - n_{e0}$ at $\omega_p t = 450$



$n_e - n_{e0}$ at $\omega_p t = 1000$



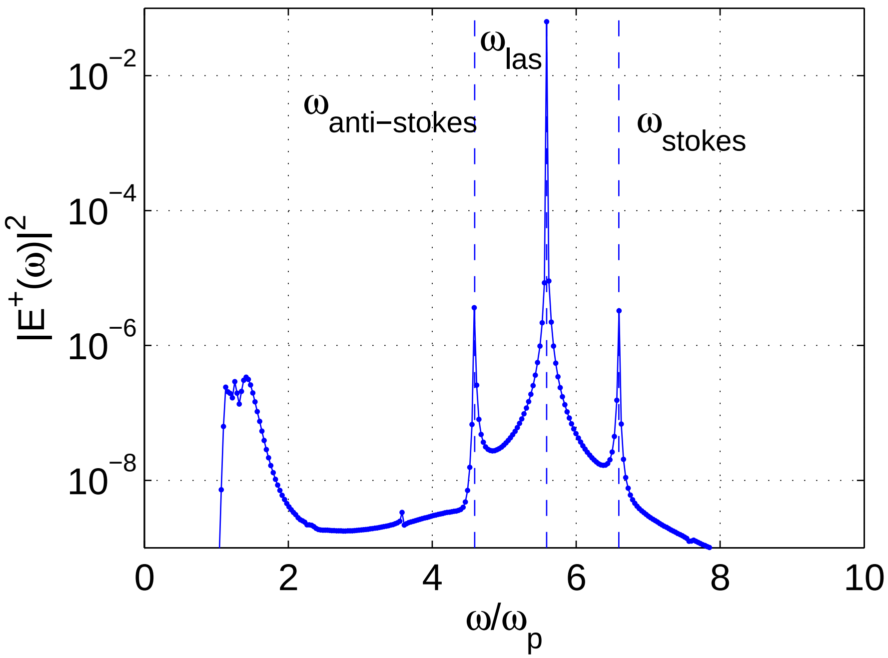
$|n_e(\omega)|^2$ at $\omega_p t = 450, x\omega_p/c = 30 - 120$



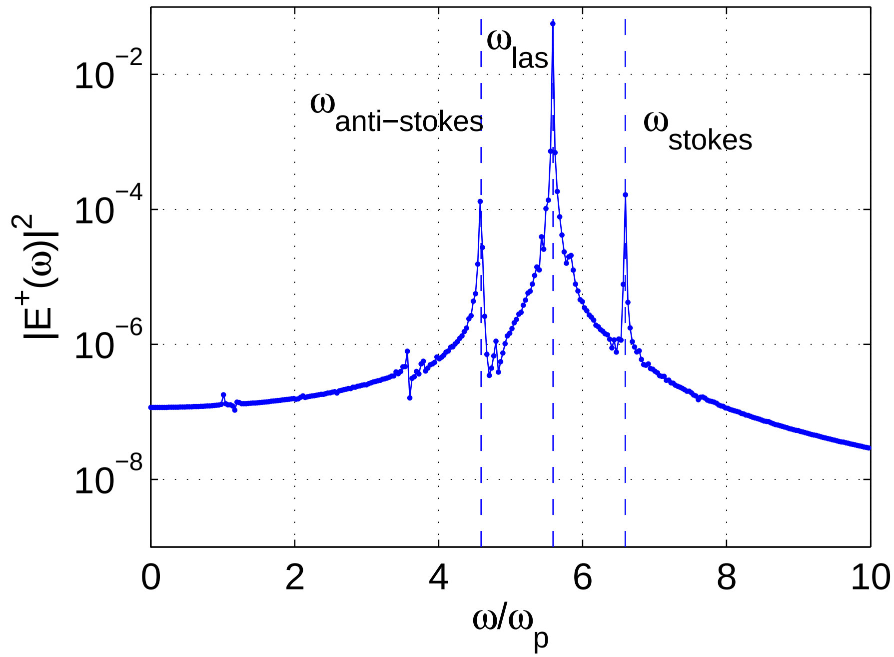
$|n_e(\omega)|^2$ at $\omega_p t = 1000, x\omega_p/c = 30 - 120$

Transmitted Light E^+ at Right Boundary

$|E^+(\omega)|^2$ at $x\omega_p/c = 160, \omega_p t = 200 - 400$



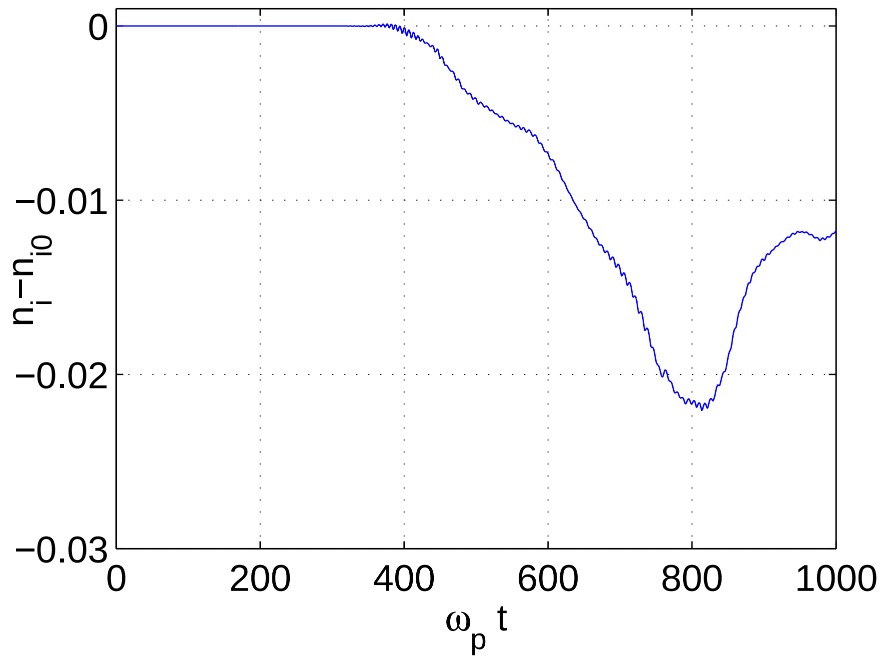
$|E^+(\omega)|^2$ at $x\omega_p/c = 160, \omega_p t = 800 - 1000$



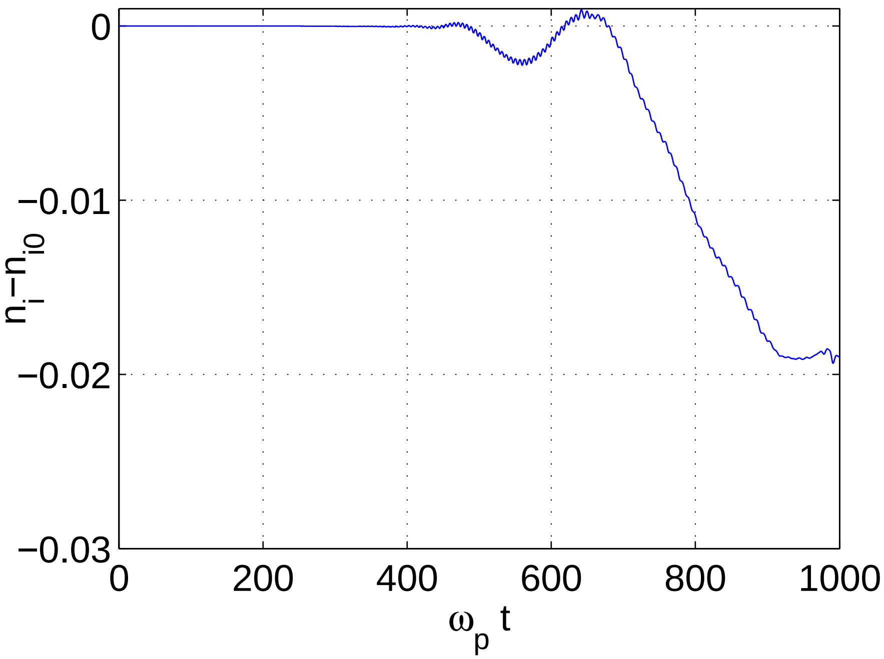
Forward Raman scattering, both Stokes and anti-Stokes, are present and grow over time

Ion Density Fluctuations vs. t

$n_i - n_{i0}$ at $x\omega_p/c = 40$



$n_i - n_{i0}$ at $x\omega_p/c = 120$



Run too short for ion acoustic wave to go through several cycles, e.g.:

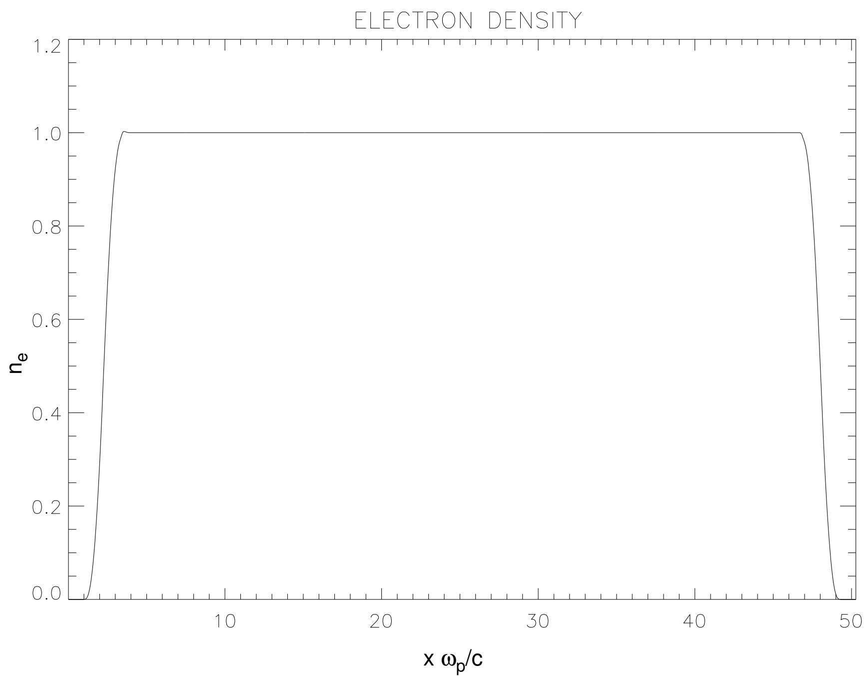
$$c_a = 8.3 \cdot 10^{-4} \qquad k_{IAW-SBS} = 11 \quad \rightarrow \quad \tau = \frac{2\pi}{c_a k_{SBS}} = 688$$

Future Work: Parallelize, $1\frac{1}{2}$ -D Code

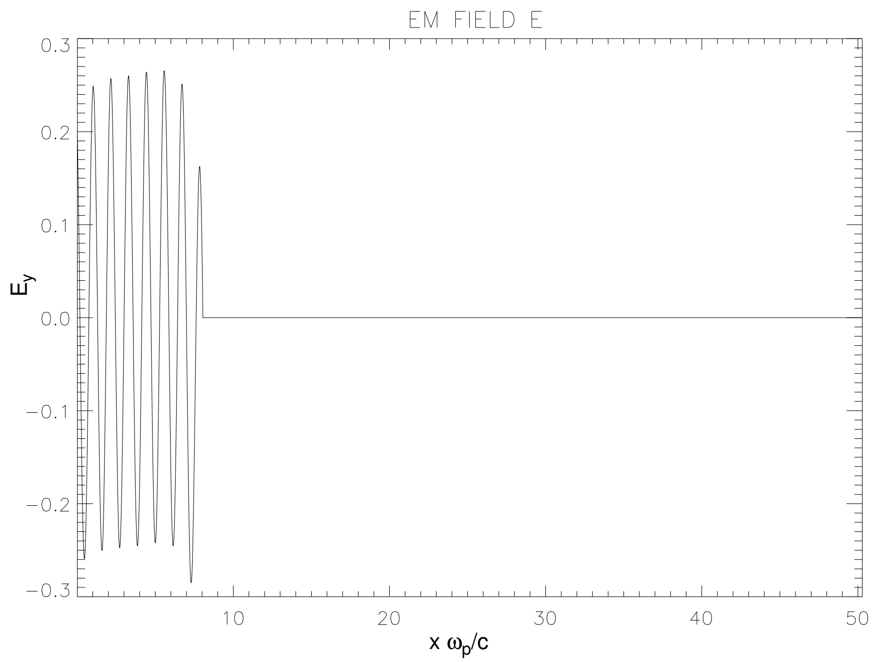
Preliminary results from $1\frac{1}{2}$ -D code that treats $v_y \ll c$ kinetically: $f_s(x, p_x, v_y, t)$.

$$\frac{\partial f_s}{\partial t} + v_x \frac{\partial f_s}{\partial x} + q_s(E_x + v_y B_z)\frac{\partial f_s}{\partial p_x} + \frac{q_s}{\gamma m_s}(E_y - v_x B_z)\frac{\partial f_s}{\partial v_y} = 0$$

n_e at $\omega_p t = 8$



E_y at $\omega_p t = 8$



Comments? Reprint? Please leave address.